

Dec. 2013, ULB

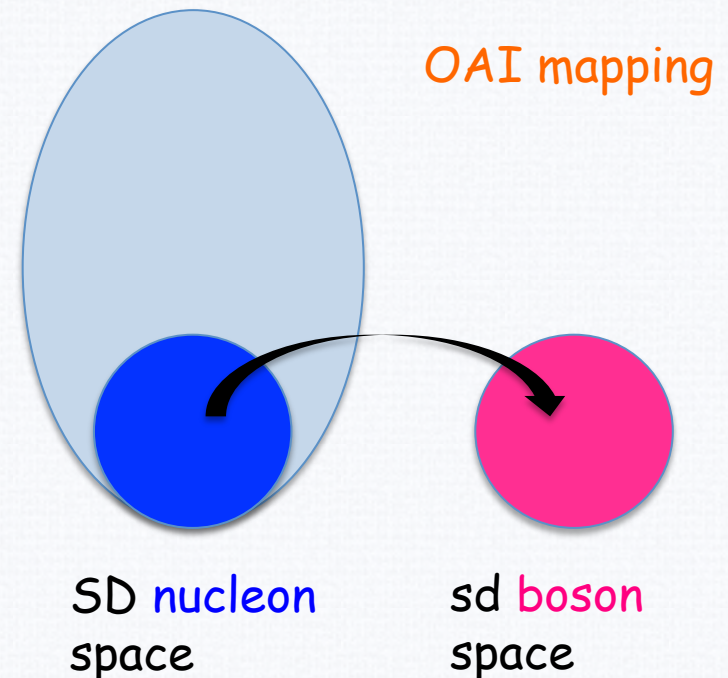
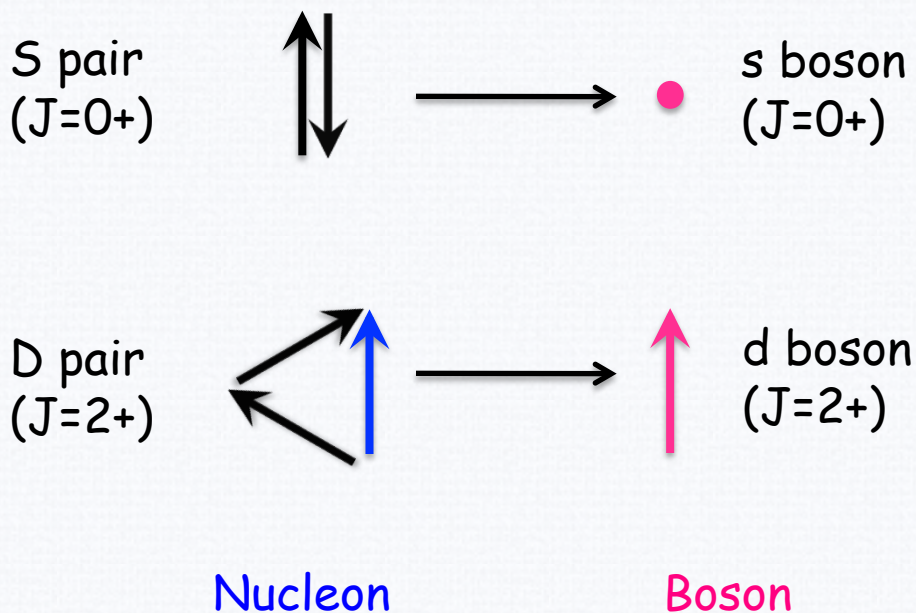
Microscopic IBM for shape coexistence

K. Nomura (GANIL)

IBM and the “microscopic foundation”

- Ingredients: collective pairs of valence nucleons
- Derivation from **shell model**. Good for modest deformation (near spherical and γ -soft systems)
- **General cases ?**

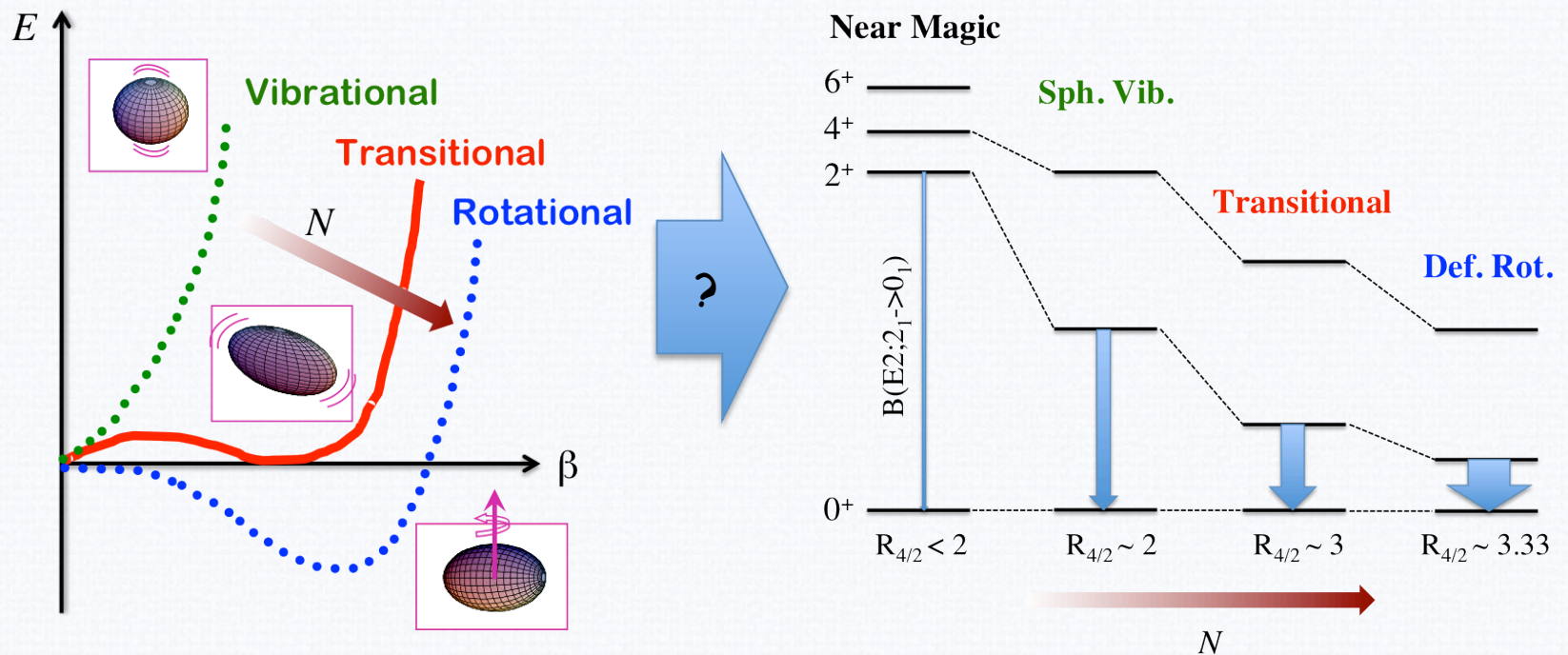
- A. Arima and F. Iachello (1974)
- T. Otsuka, A. Arima, F. Iachello (1978)
- T. Mizusaki, T. Otsuka (1997)



This work - Basic idea

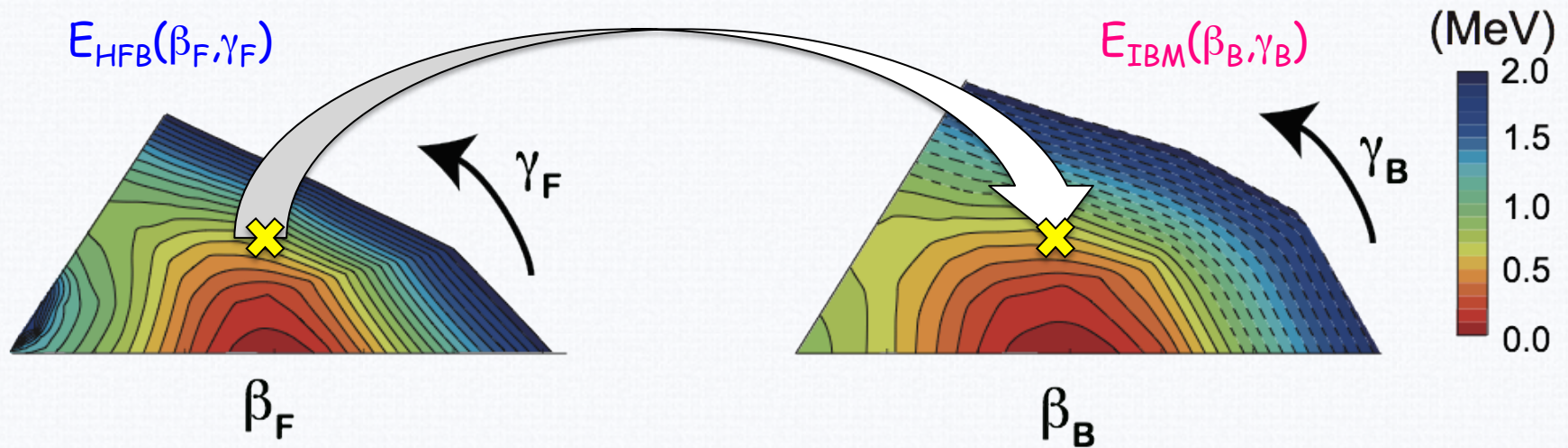
(Potential) energy surface of **mean field model** (Skyrme, Gogny, RMF, etc) can be a good starting point for intrinsic properties, directly related to shape.

We construct an IBM(-2) Hamiltonian by the **mapping** from **mean-field model**.



"Mapping" mean field

K.N., N. Shimizu, T. Otsuka,
PRL101, 142501 (2008)



Total energy for **nucleon**,
from the constrained HF(B)

Total energy for a **boson**
condensation (intrinsic state)

An IBM Hamiltonian is determined by the approximate
equality: $E_{\text{HFB}}(\beta_F, \gamma_F) \sim E_{\text{IBM}}(\beta_B, \gamma_B)$
→ energy levels and wave functions with good J, N, P, ...

PES fit by Wavelet

- extracts global topology of PES: curvature, minimum, etc.
- fixes IBM parameters unambiguously

$$\tilde{E}(\delta\beta, \beta) = \frac{1}{\sqrt{\delta\beta}} \int E(\beta') \Phi^* \left(\frac{\beta - \beta'}{\delta\beta} \right) d\beta'$$

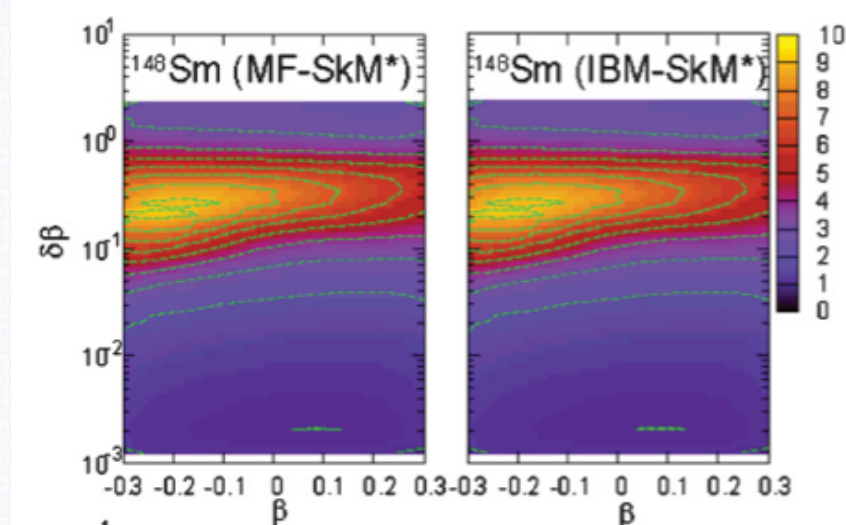
Diagram illustrating the wavelet transform fit of the Potential Energy Surface (PES):

- scale (frequency)** (red text) points to $\delta\beta$ in the denominator of the wavelet function.
- coordinate** (blue text) points to β in the wavelet function.
- PES** (blue text) points to $E(\beta')$ in the integrand.
- basis (wavelet)** (green text) points to Φ^* in the integrand.

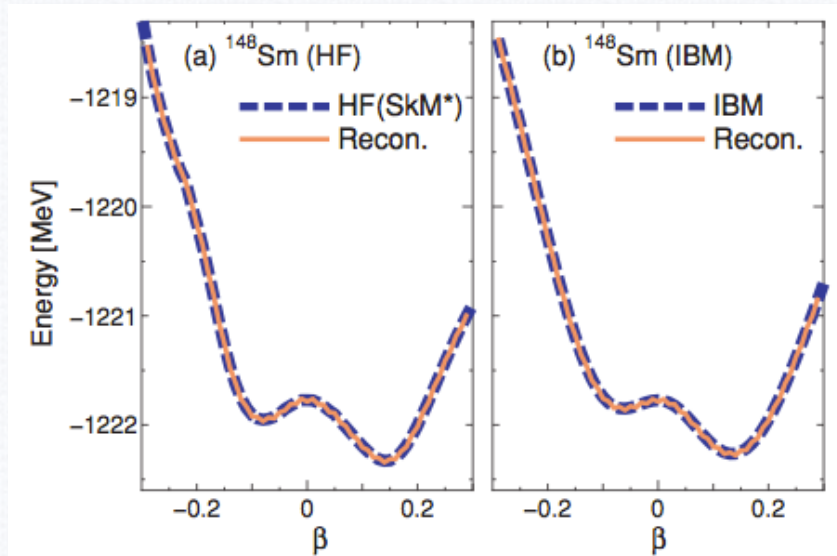
General rules:

- topology around minimum should be reproduced
- do not try to fit $\beta \gg \beta_{\min}$

χ^2 fit for wavelet transforms



comparison in real space

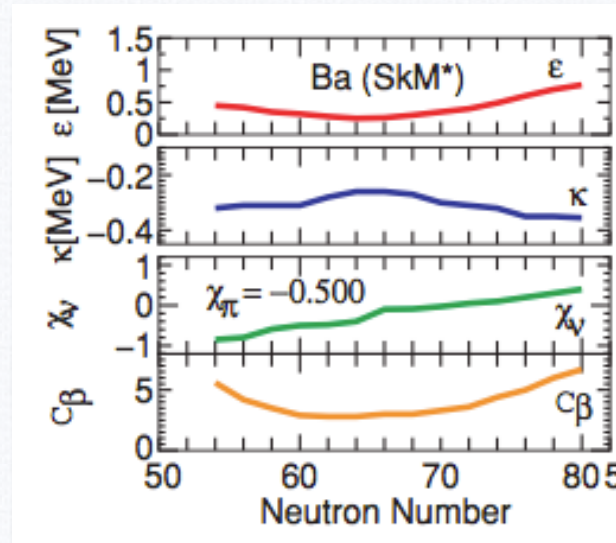
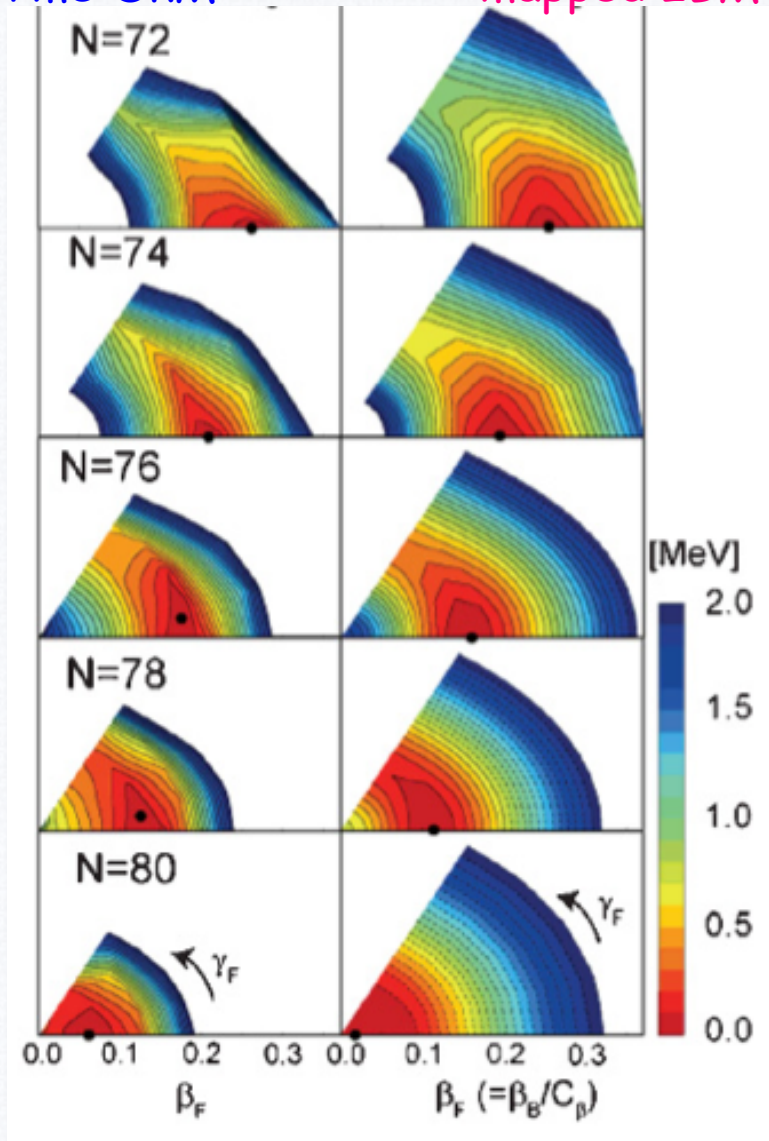


Example: Ba isotopes

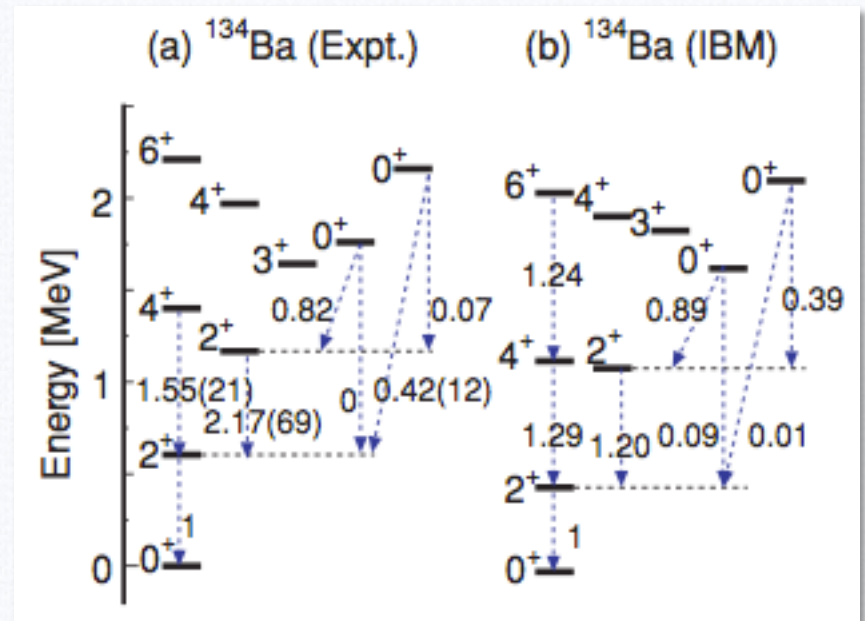
K.N., N. Shimizu, T. Otsuka,
PRC81, 044307 (2010)

Skyrme-SkM*

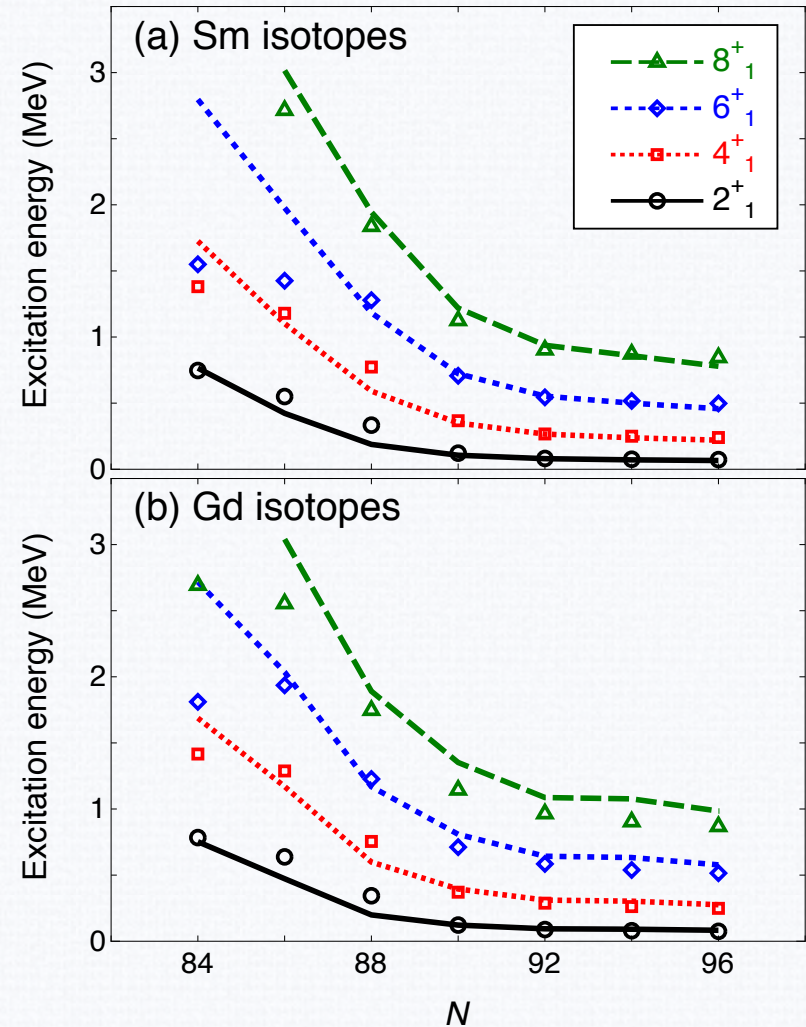
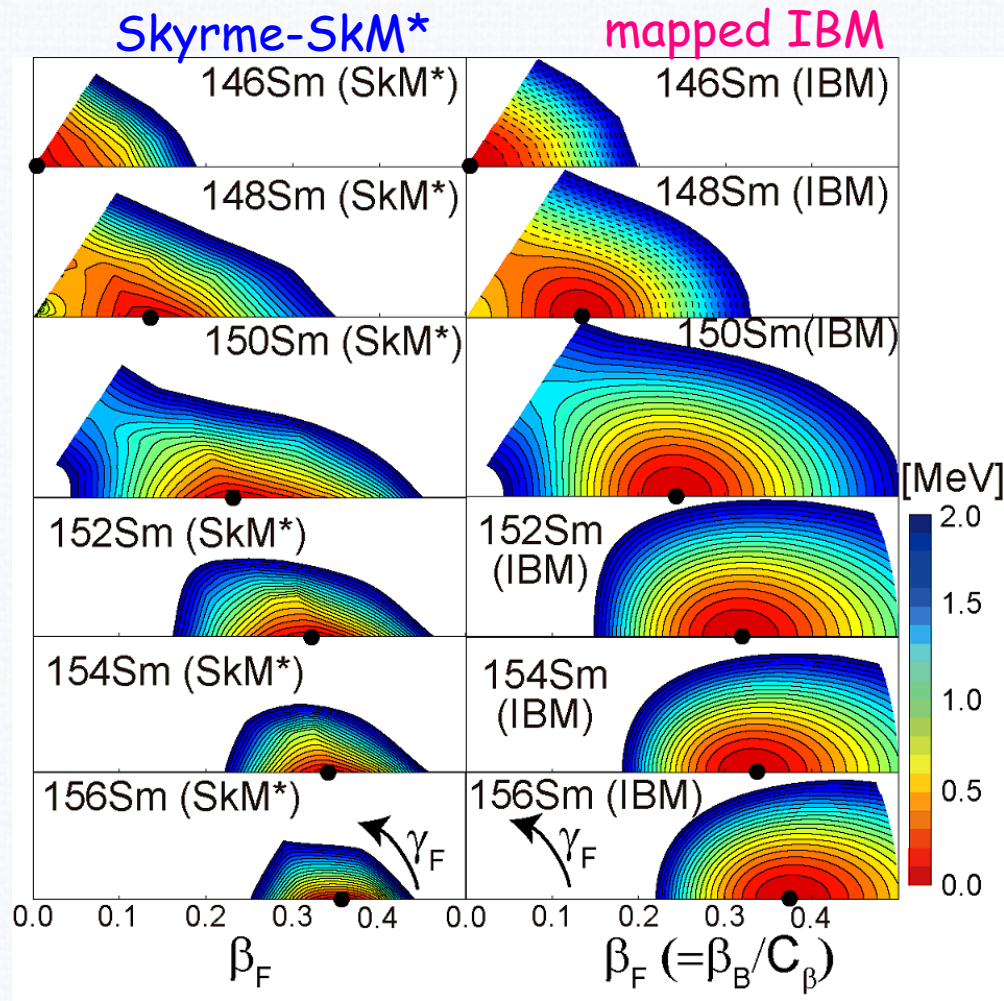
mapped IBM



Derived
parameters



How about well deformed cases?



K.N., T. Otsuka, N. Shimizu et al., PRC83, 041302(R) (2011)

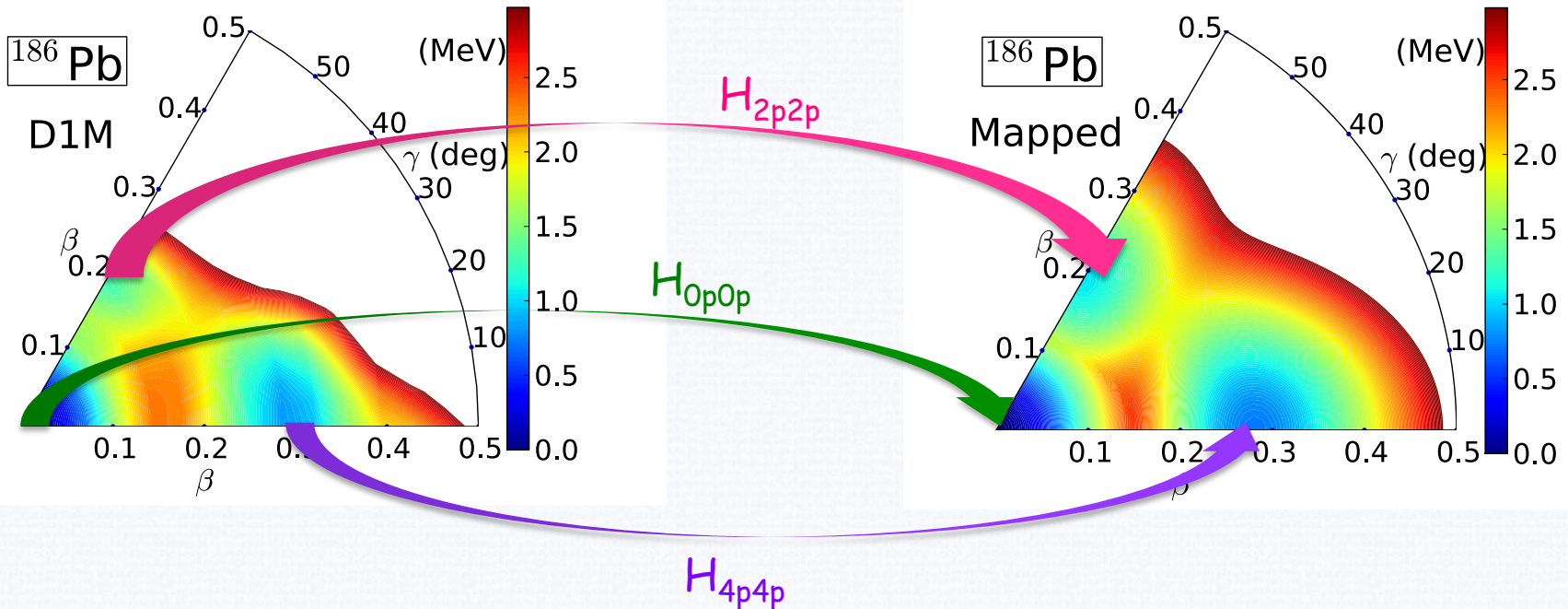
Configuration mixing in IBM

Consider mixing of different Hamiltonians for ph excitations, each associated to a mean-field minimum

$$H = H_{0p0h} + (H_{2p2h} + \Delta_{2p2h}) + (H_{4p4h} + \Delta_{4p4h}) + V_{02} + V_{24}$$

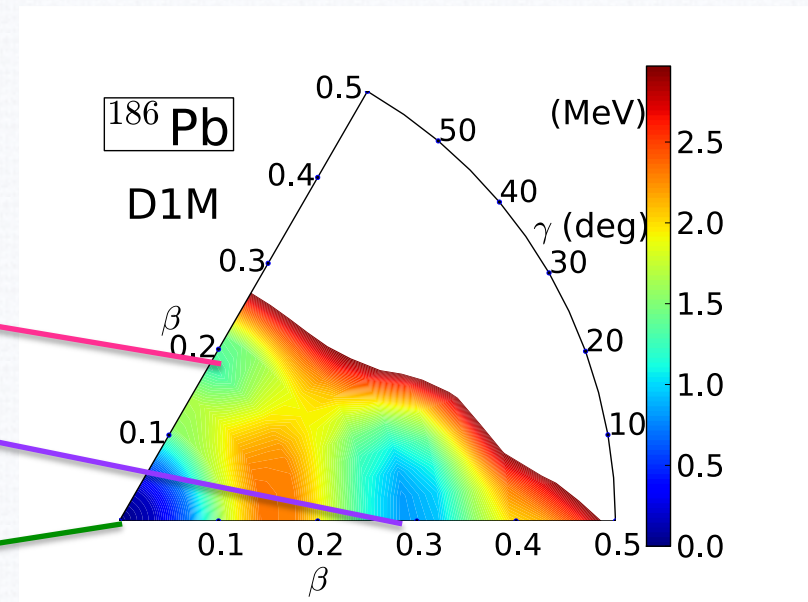
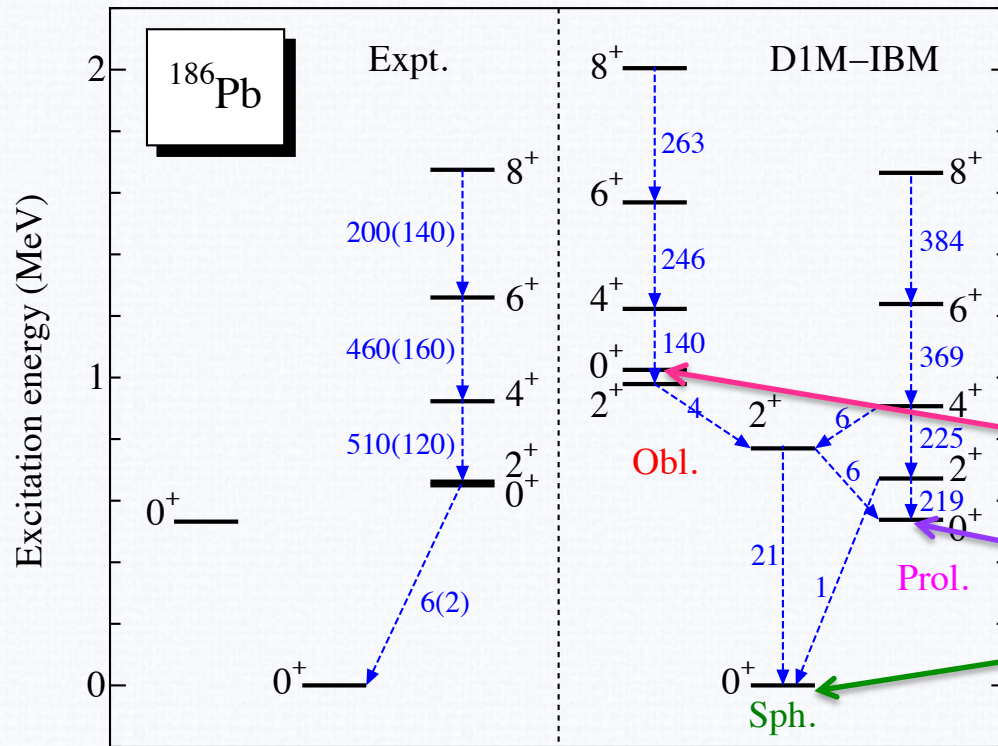
energy difference between minima

barrier height



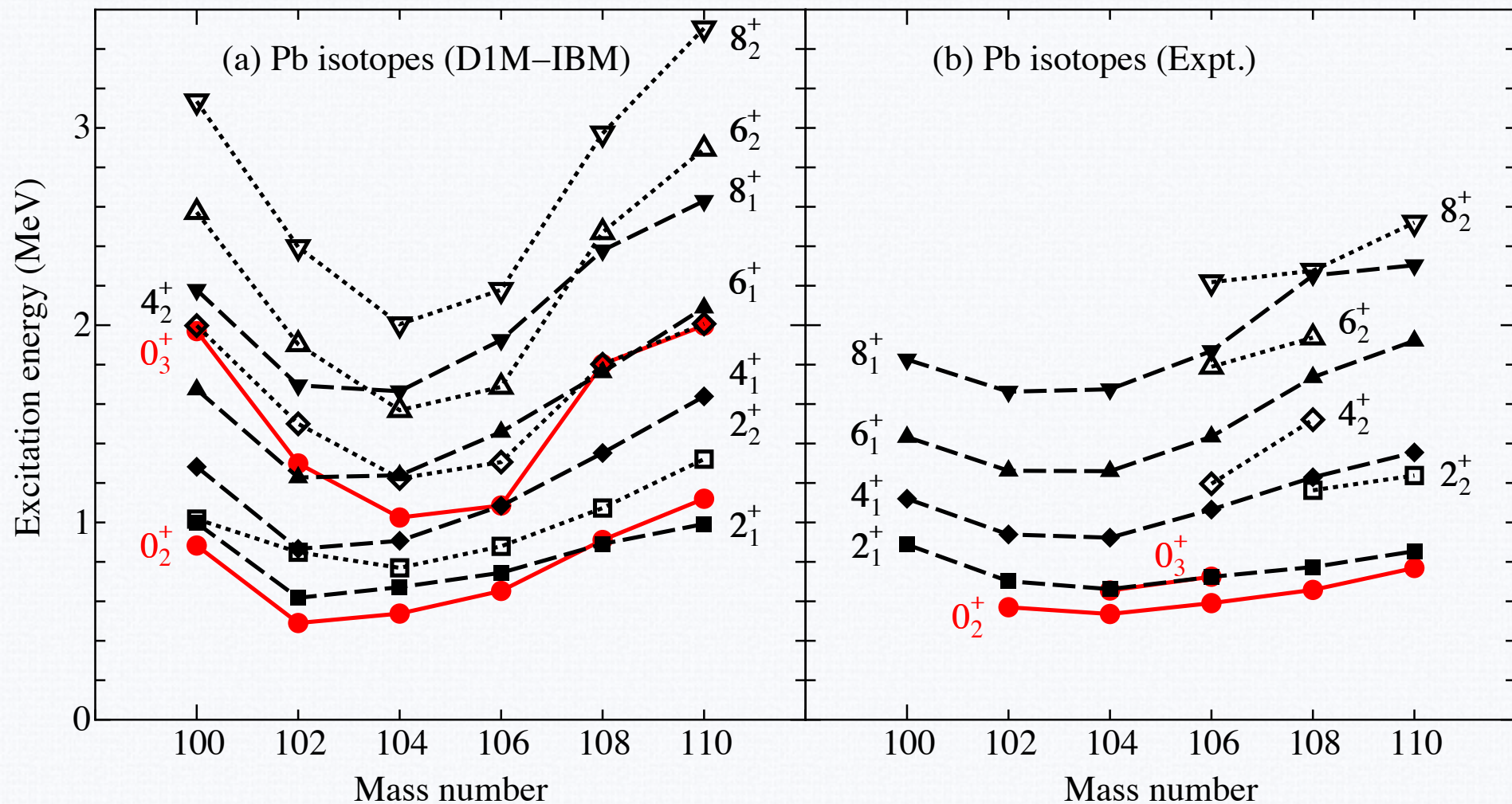
Level scheme: ^{186}Pb

from Gogny-D1M



K.N., R. Rodriguez-Guzman, L. M. Robledo et al., PRC86, 034322 (2012)

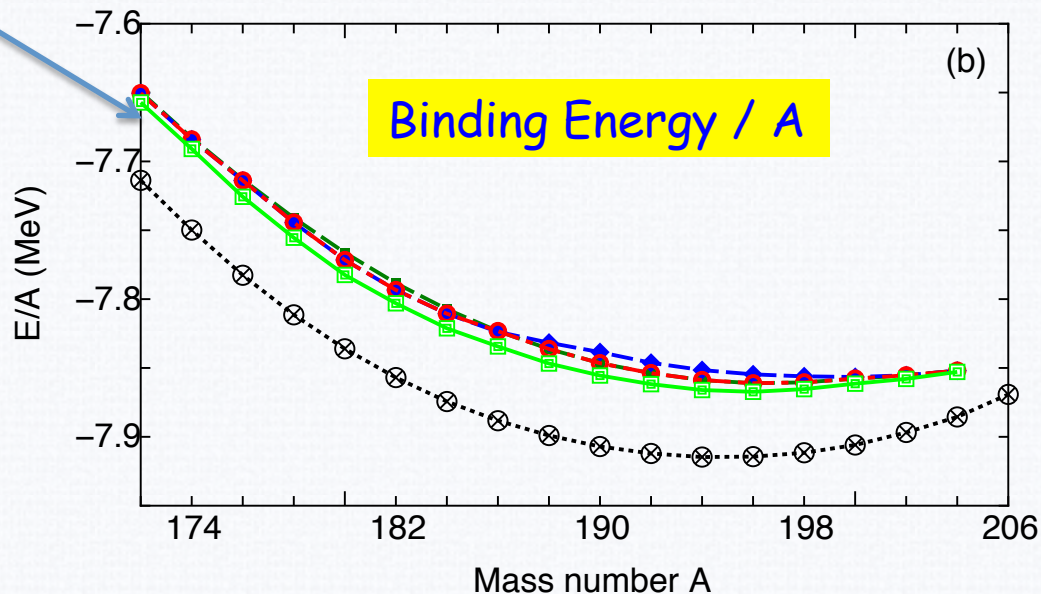
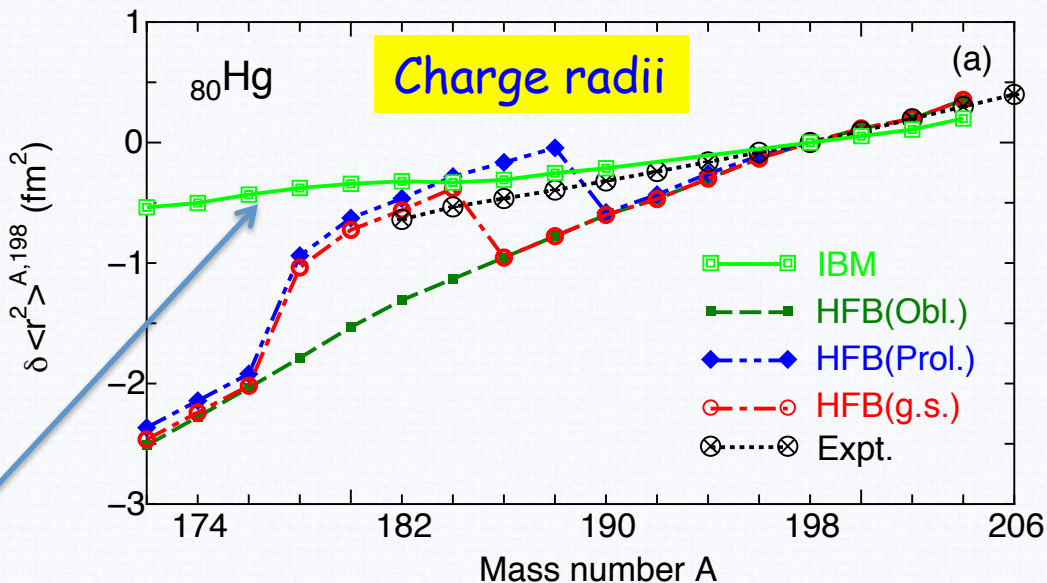
Energy-level systematics in Lead



K.N., R. Rodriguez-Guzman, L. M. Robledo et al., PRC86, 034322 (2012)

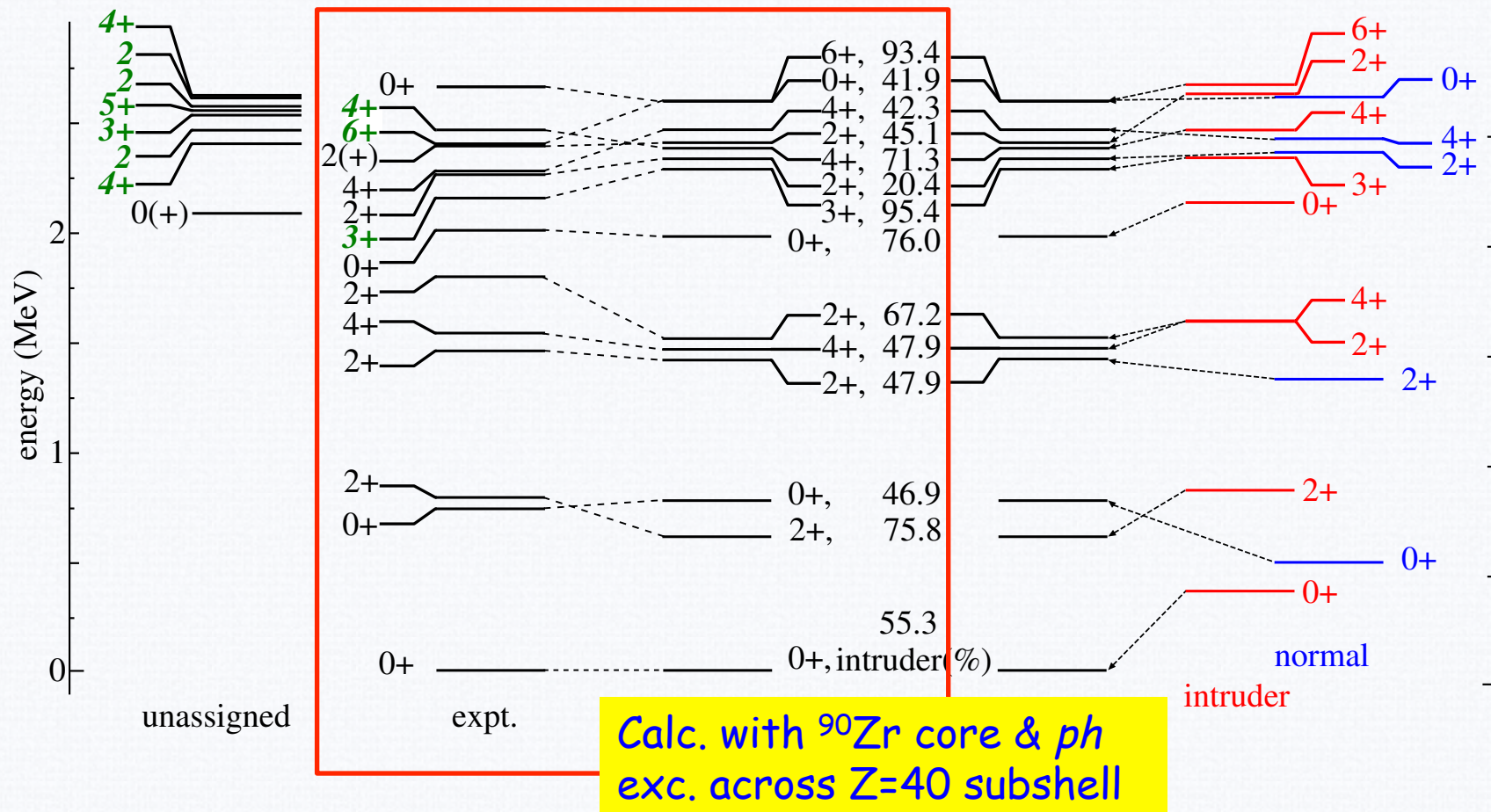
Ground-state properties

from diagonalization of
mapped IBM Hamiltonian



Coexistence in $A \sim 100$: ^{98}Mo

T. Thomas, K.N., V. Werner et al.,
PRC88, 044305 (2013)



- Coupling of **spherical normal** and **γ -soft intruder** configurations
- Large mixing in the ground state

Coexistence in $A \sim 100$: ^{98}Mo

T. Thomas, K.N., V. Werner et al.,
PRC88, 044305 (2013)

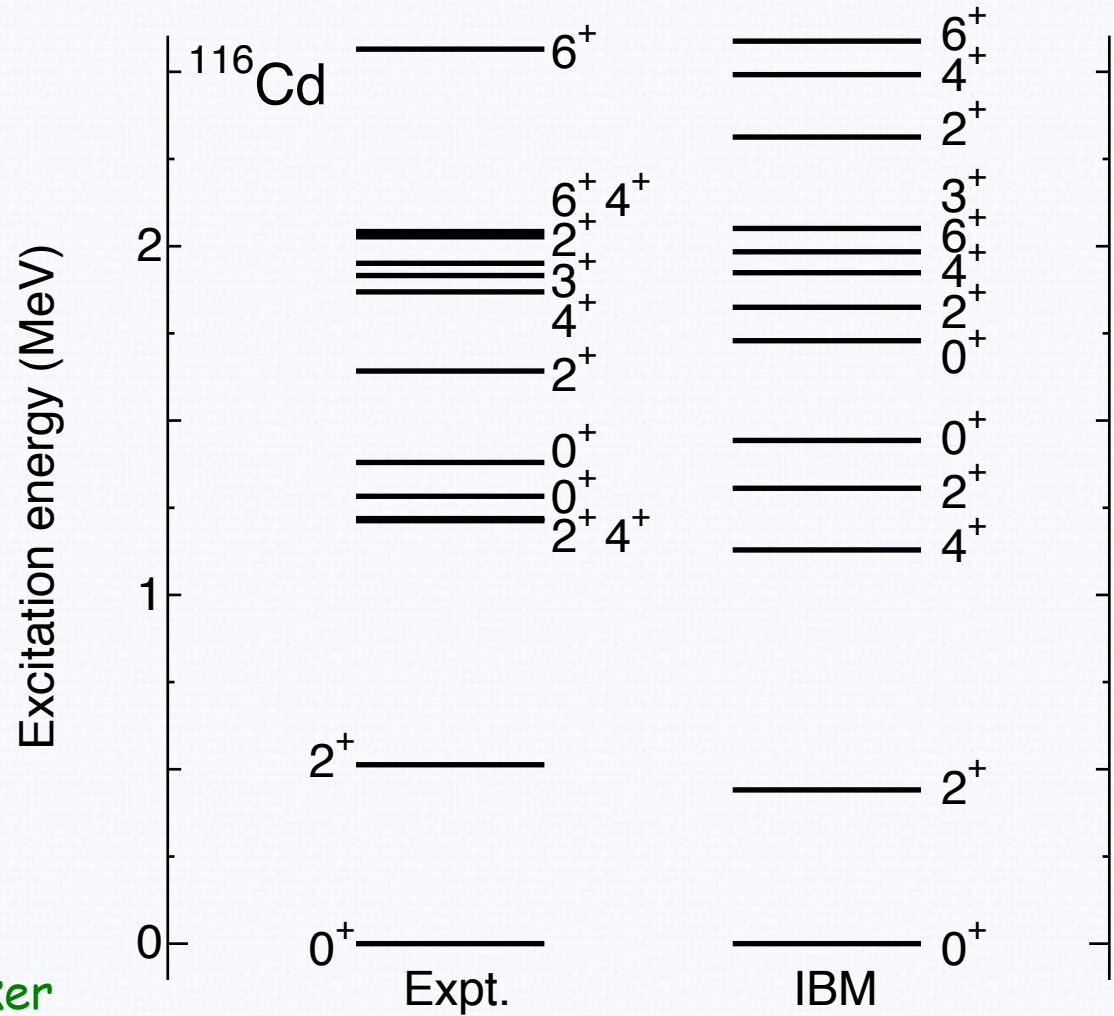
B(E2) (in W.u.)						
E_{level} (keV)	J_I^π	E_γ (keV)	J_F^π	$B(E2)_{\text{theor}}$	$B(E2)_{\text{expt}}$	δ_{expt}
787.26 ^a	2_1^+	787.26	0_1^+	27	21.4^{+11}_{-10}	
		52.6	0_2^+	256	$280 (40)^b$	
1432.29 ^a	2_2^+	644.70	2_1^+	22	47.8^{+132}_{-100}	
		697.10	0_2^+	8	2.5^{+8}_{-6}	
		1432.29	0_1^+	0.03	1.0^{+2}_{-1}	
1509.74 ^a	4_1^+	722.48	2_1^+	49	$49.1^{+5.5}_{-4.5}$	
1758.32 ^a	2_3^+	326.05	2_2^+	13	4.7^{+189}_{-23}	$-0.17 (22)$
		971.03	2_1^+	6	3.2^{+134}_{-16}	$-0.97 (14)$
		1023.61	0_2^+	7	7.8^{+286}_{-34}	
2206.74	2_4^+	1419.48	2_1^+	1.3	$1.7 (2)$	$-0.33 (11)$
2333.03	$2_5^{(+)}$	900.85	2_2^+	1	1.6^{+8}_{-4}	$-0.15^{+0.19}_{-0.20}$
2343.26 ^c	6_1^+	833.52	4_1^+	56	$10.1 (4)$	

Strong configuration mixing evident from the enhanced $2_1 \rightarrow 0_2$ and $2_2 \rightarrow 2_1$ E2 transitions

More complicated case: Cd isotopes

- from Skyrme SLy6
- ^{132}Sn core with proton excitation across $Z=50$

with J. Jolie, P. Van Isacker
et al., in progress



Summary

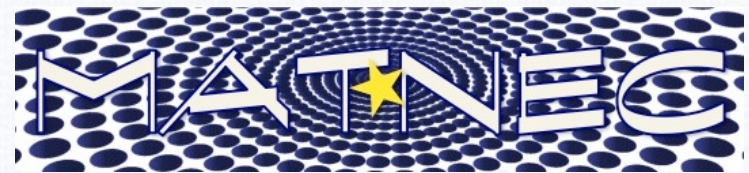
New formulation of the IBM has been developed. Bridge over the gap between IBM and mean-field model.

- Provides physical observables in **lab frame**
- Valid for **general cases**:
 - Sph. vib., def. rot. and transitional nuclei
 - Triaxial system (not mentioned)
 - **Shape coexistence and configuration mixing**
- Prediction is possible on **heavy exotic nuclei**.

Acknowledgement

My collaborators: T. Otsuka (U. Tokyo)
N. Shimizu (U. Tokyo)
L. M. Robledo (Madrid)
R. Rodriguez-Guzman (Rice U)

European Commission for the financial aid



Thank you for your attention!

Question for deformed nuclei

Physica Scripta, Vol. 22, 468—474, 1980

Features of Nuclear Deformations Produced by the Alignment of Individual Particles or Pairs

Aage Bohr and Ben R. Mottelson

The Niels Bohr Institute and Nordita, Blegdamsvej 15, DK-2100 Copenhagen, Denmark

Received May 8, 1980

by Bohr and Mottelson, "SD truncation for IBM may not be sufficient for deformed rotational system"

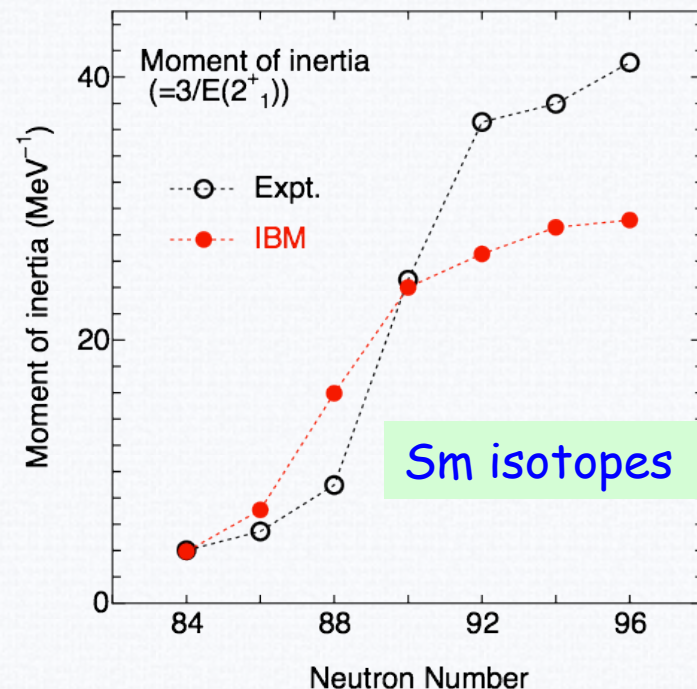
Many debates over the validity of IBM. Unified description?

- Nilsson-BCS model (T. Otsuka et al. 1982; D. Bes et al. 1982)
- renormalization of G pair or introducing g boson (T. Otsuka & J. N. Ginocchio 1985; T. Otsuka & M. Sugita, 1988)
- conventional boson mapping (M. R. Zirnbauer 1984, etc)
- J projection on intrinsic state (N. Yoshinaga et al. 1984)
- ...

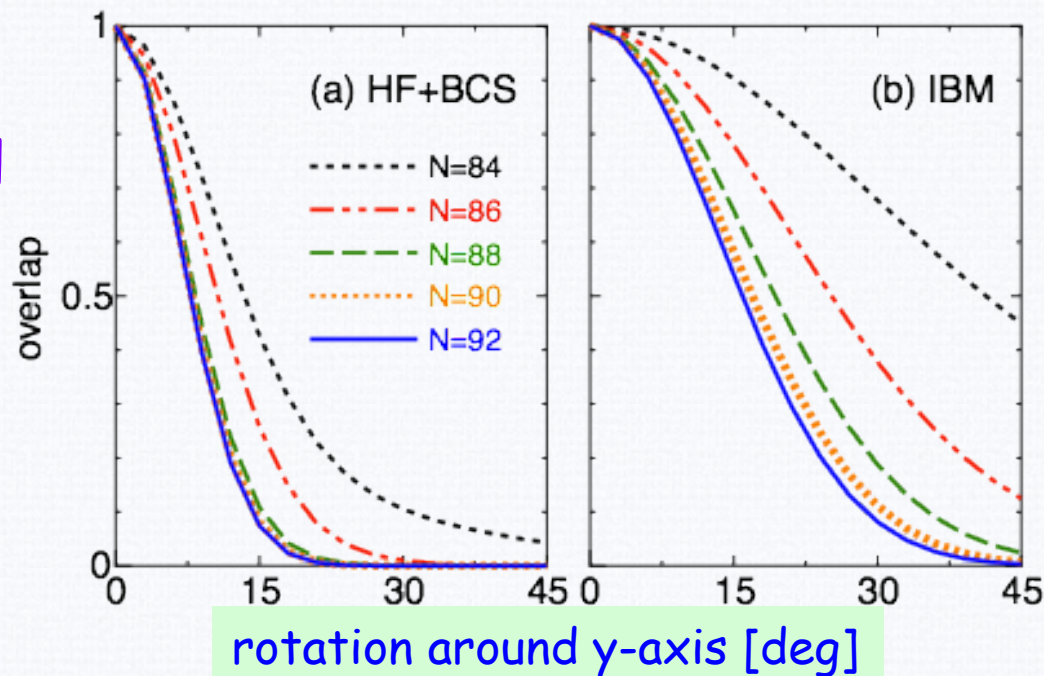
Problem with deformed rotor

Big discrepancy in the **rotational band (moment of inertia)** of strong axially deformed nuclei.

[cf. question by Bohr & Mottelson]



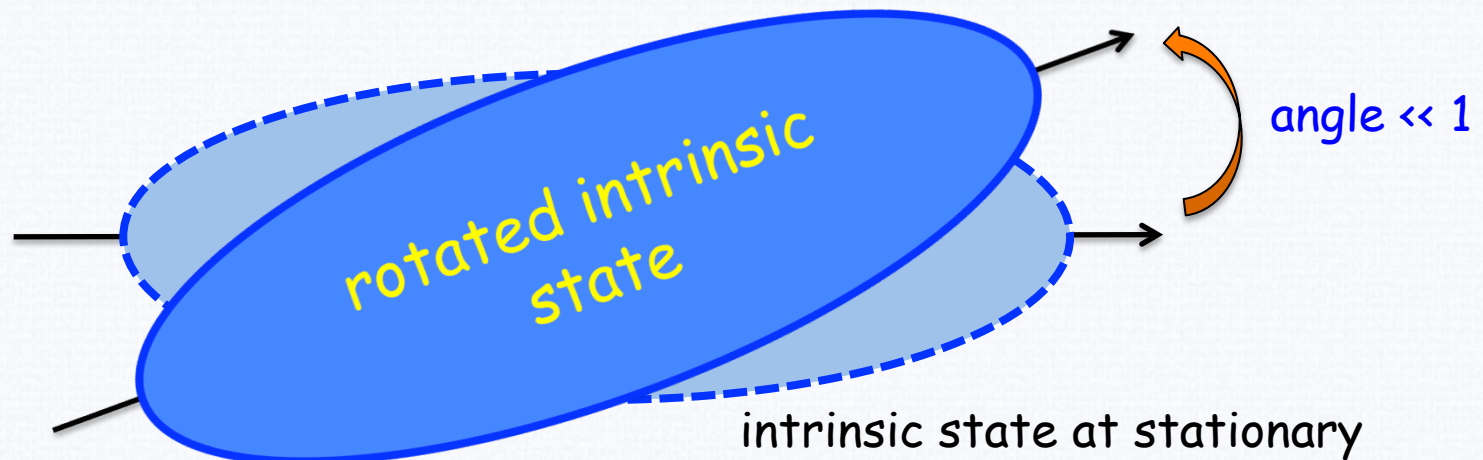
This can originate from the difference in the **rotational property of the intrinsic wave function** between fermion and boson.



Mapping “rotational response”

Basic property related to rotation should be reproduced by bosons. We consider the rotational energy (rotational response).

Some interaction term which changes energy but not change the wave function can be included. LL term.



Ref: K.N., T. Otsuka, N. Shimizu, and L. Guo, PRC83, 041302(R) (2011)

Boson Hamiltonian to be fixed

$$\hat{H}^B = \underbrace{\epsilon(\hat{n}_{d\pi} + \hat{n}_{d\nu}) + \kappa \hat{Q}_\pi \cdot \hat{Q}_\nu}_{\text{Main part}} + \underbrace{\kappa' \hat{L} \cdot \hat{L}}_{\text{LL term}}$$

1) Main part \leftarrow Mapping of the PES (static case)

$$\langle q^F | \hat{H}^F | q^F \rangle \sim \langle q^B | \hat{H}^B | q^B \rangle$$

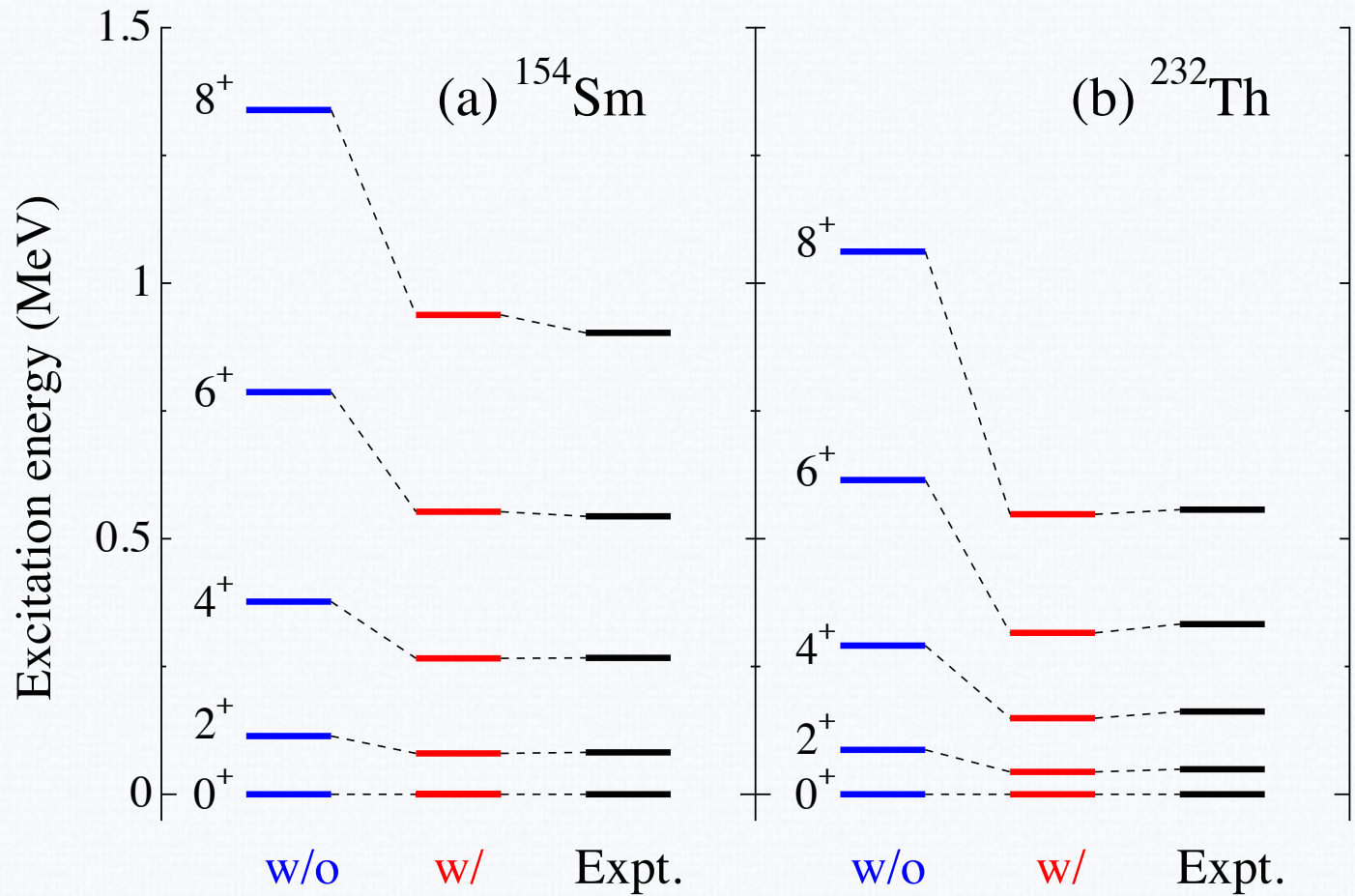
2) LL term \leftarrow rotational response at a particular deformation, e.g., minimum (dynamical case). Main part is kept in this step.

$$\delta \langle q^F | \hat{H}^F e^{-i\delta\theta J_y} | q^F \rangle \sim \delta \langle q^B | \hat{H}^B e^{-i\delta\theta J_y} | q^B \rangle$$

angle $\delta\theta \ll 1$

Cranking mom. of inertia is equated between nucleons and bosons.

Importance of LL term for rotational band

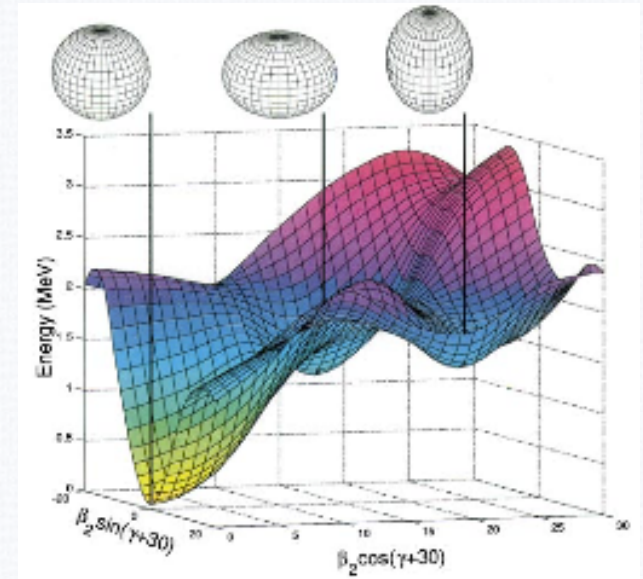


No adjustment to the experimental data.

Nuclear Structure around ^{208}Pb

Richness in nuclear shape phenomena:

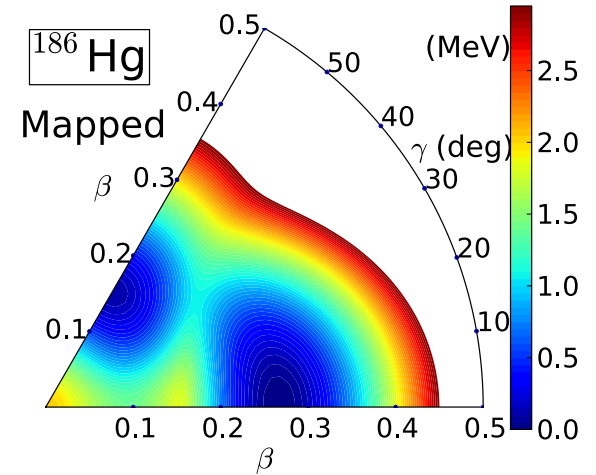
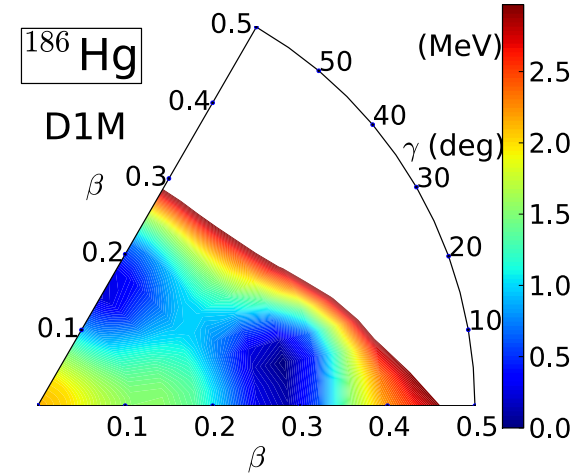
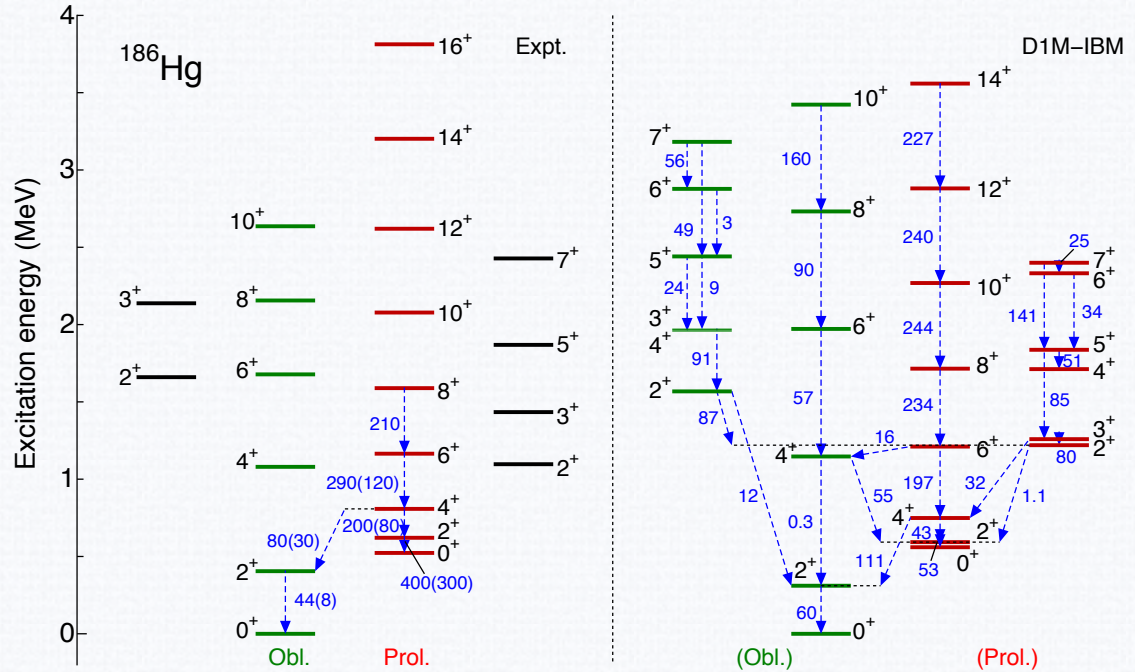
- Competing shapes near the ground state, e.g., ^{186}Pb
- Prolate-oblate transition, soft nuclear shapes
- Superdeformation
- ...



Theory:

(beyond) mean field, phenomenological IBM, ... etc.

- Configuration mixing in IBM (Duval & B. R. Barrett, '82)
- Nilsson-Strutinsky method (W. Nazarewicz, '93)
- Skyrme+GCM (T. Duguet et al., '03; M. Bender et al., '04; J. Yao et al., '13)
- Gogny+GCM (R. Rodriguez-Guzman et al., '04)
- RMF (T. Niksic et al., '02)

$$^{186}\text{Hg}$$


K.N., R. Rodriguez-Guzman, L. M. Robledo, PRC87, 064313 (2013)